RECENT STABILITY COMPARISONS WITH THE JPL LINEAR TRAPPED ION FREQUENCY STANDARDS¹

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Abstract

The ¹⁹⁹I Ig⁴ research frequency standards I. ITS-1 and 1.ITS-2 were developed to provide continuous, reliable, high stability performance. For simplicity, a 'Wig lamp is used for state selection and a helium buffer gas for ion cooling. In a preliminar y 9 day comparison between the trapped ion standards, the Allan deviation was $\sigma_y(\tau) = 1 \times 10^{-13}/\tau^{1/2}$ and a fractional frequency stability of 6×10^{-16} measured for averaging times gi cater than 10^5 seconds. A 40 day comparison of 1.11 S-2 against an auto-tuned H-maser referenced to UTC-NIST puts an upper limit on long term drift of 1.1'1'S-? of I.?(1.4)× 10^{-16} /day.

Introduction

Trapped ion frequency standards show great promise towards fulfilling several intermediate and long term frequency and timing needs. Jon trap based standards have the main advantage that the. ion (atomic oscillator) is confined only by electromagnetic fields. Perturbations due to collisions are greatly reduced and ions can in principle be held indefinitely allowing for extremely long interrogation times. The 199Hg tion is particularly well suited for frequency standards because the large mass and: 40.5 GHz ground state hyper fine splitting reduce sensitivity to thermal and magnetic variations, Research standards 1.1'1'S-1 and 1. ITS-2 were developed to provide continuous high stability operation. These linear ion trap [1] standards (1 JTS) use a 2021Ig lamp to generate J 94 nm light for optical state selection [2] and helium buffer gas to cool the ions to near room temperature [3].

The microwave ${}^2S_{1/2}(F=0,m_F=0)$ to ${}^2S_{1/2}(F=1,m_F=0)$ hyperfine transition of ${}^{199}Hg^*$ has a measured $Q>2\times10^{12}$ [5]. Good signal to noise is achieved with as many as 3×10^7 ions in a linearion trap. Several local oscillators (1,()) have been used, including a good quartz crystal, a JI-mast.I, or the Superconducting Cavity Mascr Oscillator (SCMO) [4,5,6].

Shortter m performance of " $7x10^{14}/\tau^{1/2}$ [6] is obtained using a hydrogen mast.1 as the local oscillator. With improvements to the optics—configuration, with estimate the lamp based system is capable of $4x10^{\circ}14/11^{\circ}2$.

Measured environmental sensitivity [6] indicate that au order of magnitude improvement compared to II-maser stability is possible with regulation levels still less stringent than for masers. Because a large number of mercury ions are confined at room temperature, the, second 01 der Doppler shift is the leading perturbation that will dictate the stability floor and the system accuracy. Current frequency accuracy is shout 1 0⁻¹³, though with an ion temperature measurement accurate to 1 % [7] over all accuracies of 10⁻¹⁴ should be possible. A cryogenic, laser based 19911g4 standard is currently under development at NIST [8]. This approach uses only a few ions which limits signal to noise, but has the potential of high absolute accuracy with long averaging times. Both approaches will benefit if current research to develop anultra-violet diode laser capability is successful (see e.g. [9]), The JPL standards would achieve even better short term stability, and last] cooling would become much more practical.

In this paper, we report the first 9 day stability comparison between the JPL 1 IgH trapped ion research standards LITS-2 and JUTS-1. This comparison demonstrates stabilities well into the 10"" range for averaging times longer than 1 00,000 seconds (Fig. 1). We also report a 40 day stability comparison between 1. ITS-2 and two H-masers.

Long Term Stability and Environmental Sensitivity

The limiting long term stability depends on the frequency sensitivity of the hyperfine transition to confinement and environmental perturbations. Typical operating conditions, frequency offsets, and measured sensitivities have been previously reported [6]. The accuracy

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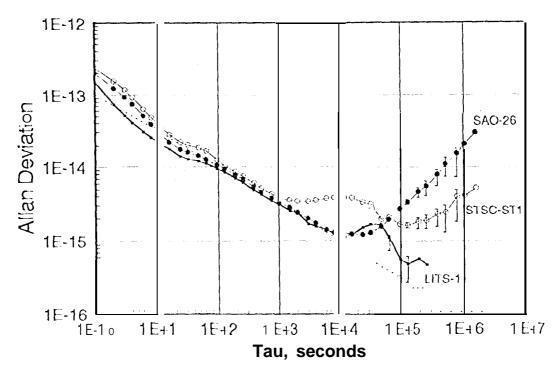


Figure 1: Forty day stability comparison of the Hg+ standard],1"1'S-2 against (a) SAO-26H-Maser, (b) STSC ST-1Auto-tuned H-Maser, and (c) a nine day comparison against the Hg+standard LJTS-1.

and limiting stability of the trappedion standards depends on how wellthese offsets are understood and held constant.

To measure the stability between the two 'Wig' standards each steers a separate VI.G-11 [1 0] }I-mast. r receiver(Fig. 2). These local oscillators consist of a good crystal oscillator phase locked to a common H-maser oscillator. Each I.O is steered at approximately 20 second intervals based on [he. error signal determined by the microprocessor controlling the ion trapinterrogation cycle. Both receivess provide a 100 MHz output and one is offset by 1 Hz. The 1 Hz beat is measured and the stability analyzed. As shown in Figure 2, the steered output of each 1.0 is also compared against other available. reference standards.

For this initial comparison both standards were operated with a 16 second microwave internogation cycle and a performance of $1 \times 10^{13} / \tau^{1/2}$. Figure 1 shows the Allan deviation of 1.1'1'S-2 compared to three separate references, 1,1'1'S-1, the H-maser SA0-26 [10], and the auto-tuned H-maser STSC-ST1 [11]. The SAO maser is useful for determining performance for averaging intervals less than 20,000. The STSC-ST1 maser has poorer short term stability, but is useful as a long, termreference[12]. The STSC-ST1 mast.r is also independently compared to UTC-NIST via GPS to provide, a reference to the international time scale.

The data shown in Figure 1 has no drift removed, though a $\sqrt{2}$ has been removed from the LITS-2 vs. LITS-1 comparison for averaging times greater than 20 seconds.

LITS-2 and LITS-1:

The stability between the, two trappedion standards reaches approximately 6x 10-'6 at 100, 000 seconds. This point consists of 6 samples and the. uncertainty is shown in figure 1. The peak atapproximately \$0,000 seconds resulted from a poor regulation circuit on 1.1"1'S-1. This is made graphically clear in Figure 3 which shows the Allan deviation of each trapped ion standard compared against SAO-26. An oscillation is observe.cl in the frequency residuals of both comparisons involving 1.1-1'S-1. The oscillation is not present in 1.1"1'S-?, which has better control electronics, in this preliminary 9 day measurement the differential drift between 1.11S-2 and LITS-1 is 3.2(2.7)x 10⁻¹⁶/day. This small drift correlates well with a known sensitivity and measured drift of the RF trapping potential of 1,1'I'S - Iduring the comparison. The long term drift of the SAO maser is measured independently by both LITS-land 1.1'1'S-2 during the same time interval of $4,4(0.3) \times 10^{-15}$ /day and $3.7(0.6) \times 1015$ /day respectively. The driftrate of the maser change.s overtime (see al so [12]). 1 for the 40 day comparison between SAO-26 and LITS-2. (Fig. 1) the measured drift is $2.4(0.3) \times 10^{-15} / \text{day}$.

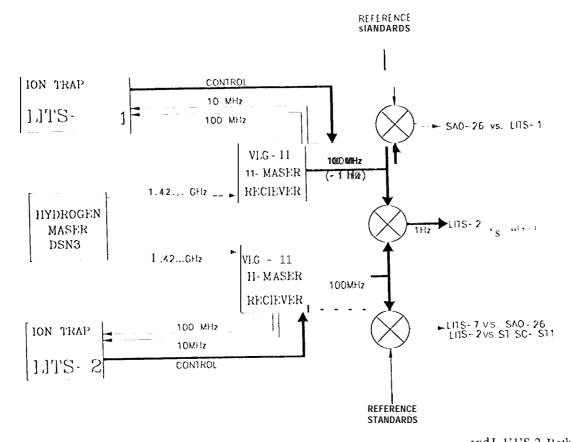


Figure 2: The measurement scheme used to compare the two trapped ion standards LATS-1 and LATS-2. Both lon traps steer a separate VLG-11 H-maser receiver.

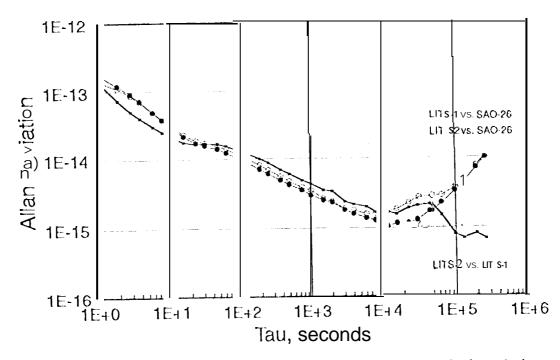


Figure 3: Nine day comparison between the Hg-trapped ion standards 1.1']'S-2 and 1.1TS-1. Both standards are also compared to the }1-maser SAO-26 over the same time interval.

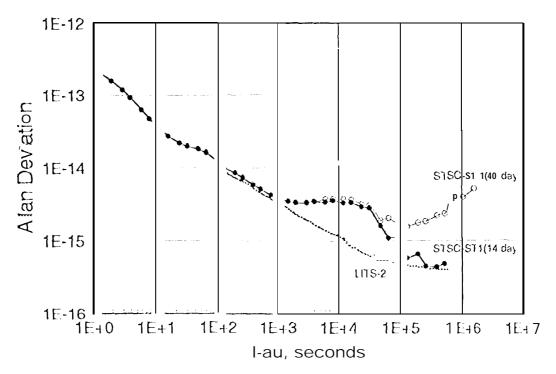


Figure 4: The Stability crf 1.1'1'S-2 compared against the H-maser STSC S'1'- 1 for a selected 14 day period between frequency jumps in the II-maser. Also shown is the complete 40 day measurement.

LITS-2 and STSC-ST1:

in Figs. 1 and 4 the Allan deviation of a 40 day comparison between LJTS-2 and the auto-tuned H-maser STSC-ST1 is shown. The differential drift between these two standards over this 40 day span is $4.7(1.6)x10^{-16}/day$. A closer examination of the time residuals shows a large 10^{-14} frequency shift 10 days into the measurement. This shift can riot be accounted for in LJTS-2 and is apparently due to a frequency jump in the STSC-ST1 maser. This frequency change is confirmed inlong term time residuals in GPS measurements comparing the STSC-ST1 H-maser to UTC-NIST. Similar frequency jumps a few times a year have also been reported elsewhere [12]. For the purpose of char.w[c~izing the.slfibility of 1.1'1'S-2, long term reference of the maser to UTC-NIST indicate, stable reference windows between frequency jumps in STSC-ST1.

Figure 4 shows the Allan deviation for a two week interval of the 40 clay comparison. During this time interval the frequency stability of the STSC-STI H-maser is particularly good as confirmed by the GPS comparison with NIST. In this 14 day comparison, the differential drift between 1.1"1'S-2 and the STSC-STI maser is measured to be 1.2(1.4)x I (1"/day. For averaging times longer than 100,000 seconds this measurement is in agreement with the performance of LITS-2 as measured by LITS-1(Fig. 1).

Electronic Improvements and Reducing Sensitivity With The Extended Linear Ion Trap (LITE)

Because of the low sensitivity to thermaland magnetic perturbations, averaging to 10^{-15} stability is accomplished with only minimal electronic control and isolation from environmental perturbations [6]. LITS-land 1.1'1'S-2 are research laboratory standards and though portable, are not highly regulated. The data presented here was obtained with the standards thermally regulated to 0.05 C and a low field differential magnetic shielding factor of only 800. The trapping potentials are run "open loop" and the ion number is not actively servoed. Several improvements to the control electronics are currently under development which should allow the standards to average with characteristic $1/\tau^{1/2}$ behavior to near 1x 10^{-12} .

In addition to relying on further electronic improvements for improved stability there are, ways to reduce fundamental sensitivity and still maintain a practical, room temperature, lamp based system. An extended version of the linear ion trap (1 JTE) is currently under development [13] which takes advantage, of the capability to easily move ions. By moving ions between two regions of a linear ion trap, the two often conflicting tasks of ion loading and optical state selection can be separated from the microwave intro-ogation region which requires an excellent magnetic

environment. Moving the ions into a long interrogation region reduces the linear ion density without sacrificing signal to noise. This not only reduces sensitivity to second m der 1 Doppler perturbations but may allow operation at lower magnetic fields.

Conclusion

A second ¹⁹⁹Hg4 tr apped ion frequency standard 1 .1'1'S-2, now under continuous operation, provides a capability for measuring stability beyond all existing frequency standards for averaging times longer than 20,000 seconds. In a recent fr equency stability comparison between ion trap standards 1 JTS-1 and 1,1"1'S-?, each standard steered as caparate VI.G-11 hydrogen maser receiver demonstrating stabilities of 6×10^{-16} for averaging limes up to 9 days. The Allan deviation of each standard was $\sigma_y(\tau) = 1 \times 10^{-13}/\tau^{1/2}$ with the differential drift measured to be $3.2(2.7) \times 10^{-16}$ /day. This remaining small drift is predominantly in 1.1'1'S-1 and correlates well with a measured drill in the trapping potential. A 40 day comparison of 1.1'1'S-? against an auto-tuned 11-maser referenced to UTC-NIST provides an upper limit on the drift of 1.1'1'S-2 of 1.2(1.4) $\times 10^{-16}$ /day.

With both standards oper s 1 in g at the previously demonstrated short term performance of $\sigma_y(\tau)$: $7x\,10^{-14}/\tau^{1/2}$ a stability of $1x\,10^{-16}$ should be possible in $5x\,10^5$ seconds given sufficient magnetic shielding and stability in the control electronics. Measured environmental sensitivities indicate this can be accomplished with regulation still less stringent than for hydrogen masers. in addition, use of a new extended linear ion trap (1.1TE) configuration should further reduce remaining sensitivity to ion number and magnetic field fluctuations, allowing for even higher stabilities.

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